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Photons' in Relativistic Plasma with Velocity Shear: A novel mechanism for power law spectra at high energies

> Mukesh Kumar Vyas Bar Ilan University, Ramat Gan, Israel Supervisor : Asaf Pe'er



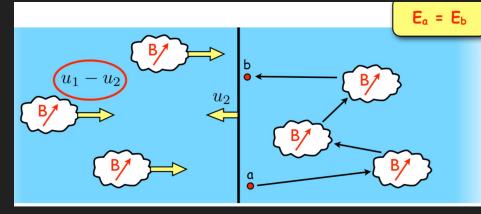


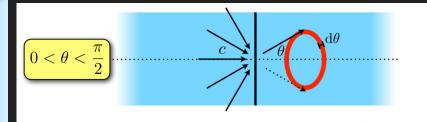


<u>Overview</u>

- Introduction : Fermi acceleration and photon energy gain in shearing flows
- **Physical mechanism** of photon energy gain
- Generation of power law spectrum at high energies
- <u>Results</u> : Comparison of estimated spectral slopes with Monte Carlo Simulations
- <u>Conclusions</u> : Significance of the work

First order Fermi acceleration: Shock acceleration of charged particles



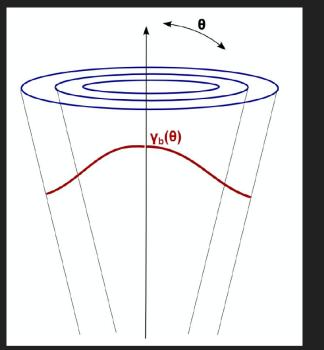


$$E' = E + \frac{E}{c} v \cos(\theta) \qquad \qquad \overbrace{E}^{\Delta E} = \frac{v}{c} \cos(\theta)$$

energy gain per half-cycle
(up->down-stream)

Electrons gain energy and produce a power law spectrum when they interact with regions having different speeds in the plasma

Particle acceleration in shearing flows



shearing flows particle acceleration

Q

About 55,700 results (0.14 sec)

An introduction to **particle acceleration** in **shearing flows** FM Rieger - Galaxies, 2019 - mdpi.com

... Shear flows are ubiquitously ... acceleration of charged particles in shearing flows and reviews some of the recent developments. Topics include the acceleration of charged particles by ... ☆ Save 切り Cite Cited by 37 Related articles All 7 versions ≫

Particle acceleration in shearing flows: efficiencies and limits

FM Rieger, P Duffy - The Astrophysical Journal Letters, 2019 - iopscience.iop.org ... We examine limits to the efficiency for **particle acceleration** in **shearing flows**, showing that relativistic **flow** speeds are required for efficient gradual **shear acceleration**. We estimate ... ☆ Save 𝔊 Cite Cited by 15 Related articles All 6 versions Web of Science: 13 ≫

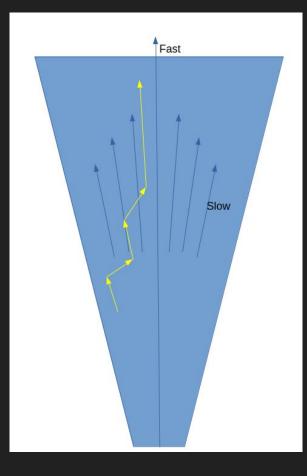
Turbulence and Particle Acceleration in Shearing Flows

FM Rieger, P Duffy - The Astrophysical Journal Letters, 2021 - iopscience.iop.org ... We explore constraints imposed by **shear**-driven instabilities on the **acceleration** of energetic **particles** in relativistic **shearing flows**. We show that **shearing** layers in large-scale AGN jets ... ☆ Save 𝔊𝔅 Cite Cited by 6 Related articles All 5 versions Web of Science: 3 ≫

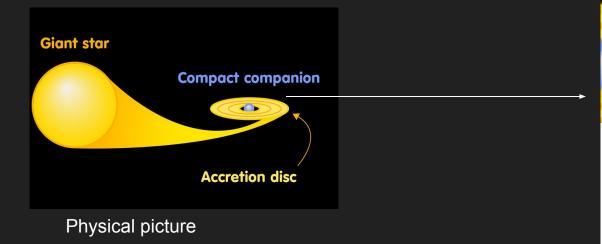
Rieger Galaxies 2019, 7(3) But what about a photon, trapped and escapes from shearing scatters?? We need investigation and a clear picture of this process.

Photon energy gain In shearing flows: the concept and motivation

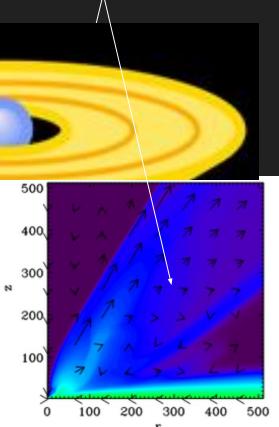
Any relativistic flow with a presence of (i) Shear, (ii) High optical depth would lead to net photon energy gain producing a power law emission of spectra



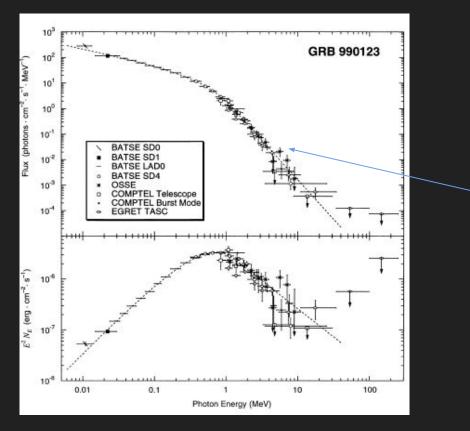
Generation of power law due to Velocity shear in the plasma



Simulations of radiation-driven winds from Keplerian discs (Raychaudhuri, Vyas & Chattopadhyay, 2021)



Attempts to explain high energy power law



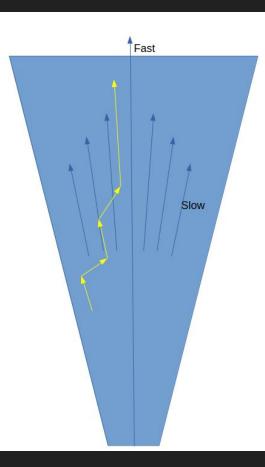
Synchrotron shock model Tavani 1996; Cohen et. el 1997; Panaitescu et. al. 1999; Frontera et al. 2000

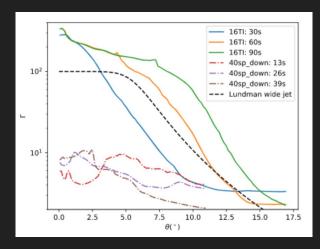
Backscattering dominated emission Eichler & Manis (2007, 08), Eichler 2004, Vyas et al. (2020, 2021)

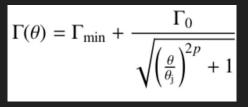
Bulk Comptonization from structured jets with photospheric emission : *Lundman et al. 2013, Vyas & Pe'er 2022 (current work*)

Piron, Comptes Rendus Physique 17, no. 6 (2016): 617-631

Photon energy gain from relativistic jets with structured Lorentz factor profile



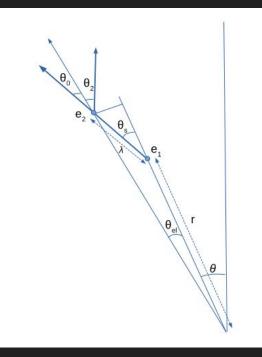


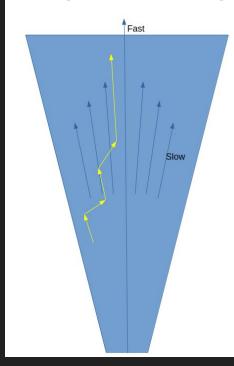


Lorentz factor (Γ) as a function of θ Zhang et al. 2003, ApJ, 586 Lundman et al. 2013, 428-3, 2430 Lundman et. al. 2014, MNRAS **440** 3292

Tyler Parsotan et al 2020 ApJ 896 139

Case of GRB jets : Energy gain in scattering process 1. Analytic estimates of spectral slopes





$$g = \frac{\varepsilon_2}{\varepsilon} = \frac{1 - \beta_2 \cos(\theta_s - \theta_{el})}{1 - \beta_2 \cos(\theta_2)}$$

$$\frac{\Gamma_2}{\Gamma} = 1 + \frac{\delta\Gamma}{\Gamma} = 1 + \sum \frac{\partial \log\Gamma}{\partial x^i} \delta x^i$$

Energy gain in scattering process

If the local mean free path measured in the lab frame is $\lambda(r, \theta)$, then

$$\tan \theta_{\rm el} = \frac{\lambda \sin \theta_{\rm s}}{r + \lambda \cos \theta_{\rm s}} = \frac{a \theta_{\rm s}}{1 + a \cos \theta_{\rm s}}$$

Taylor expansion allows us to estimate the gain as:

$$g(r,\theta) \approx \frac{1}{2} \left[1 + \left[1 + \sum \frac{\partial \log \Gamma}{\partial x^i} \delta x^i \right]^2 \frac{1}{(1+a)^2} \right]$$

$$\begin{array}{c} \theta_{0} \\ \theta_{2} \\ e_{2} \\ e_{2} \\ \theta_{3} \\ \theta_{3} \\ \theta_{0} \\$$

$$a \equiv \lambda/r$$

Averaged gain over all scatterings

The expectation value of the photon energy gain in the plasma is evaluated by integrating the average gain over the entire region of scattering

$$\bar{g} = \frac{1}{V} \int_{V} g_{a}(\theta, r) dV.$$

Escape probabilities for the photon from accelerating region

 $P_{\rm e}(r,\theta) = \exp[-\tau(r,\theta)]$ Escape probability at location *r*, θ



Photospheric optical depth

 $\tau_2 = \int_0^{s_0} \Gamma(1 - v \cos \theta_{\rm s}) n_{\rm e}' \sigma ds.$

Angular optical depth

$$\tau_2 = \frac{2r_{ph}}{\beta r\cos\theta_s} \left[1 - \frac{1}{(\theta - \theta_j)\gamma\cos\theta_s + 1} \right]$$

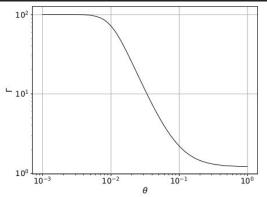


Figure 2. Lorentz factor (Γ) profile of the jet characterized by equation 13 with parameters p = 2.0, $\theta_j = 0.01$ rad, $\Gamma_0 = 100$ and $\Gamma_{\min} = 1.2$. $\theta_e = \theta_j \Gamma_0^{1/p}$. The inner jet region is for $\theta < \theta_j$ while outer region extends beyond θ_e . The region bounded within $\theta_j - \theta_e$ harbours an effective velocity shear leading to photon energy gain.

 $\tau = \min(\tau_1, \tau_2)$

$$P(r, \theta) = 1 - \exp[-\tau(r, \theta)]$$

Probability for next scattering at location r, θ without escape

Averaged probabilities from the accelerating region

The average probability of the photon to have a scattering without escape within the jet is the probability $P(r, \theta)$ averaged over the available volume V of scattering where velocity shear is present, i.e.,

$$\bar{P} = \frac{1}{V} \int_{V} P(r,\theta) dV.$$

Generation of power law spectrum at high energies

After scattering k times, there are N photons left out of N_0 photons in the accelerating region while $N_0 - N$ have escaped,

$$\frac{N}{N_0} = \bar{P}^k \text{ or } \ln \frac{N}{N_0} = k \ln \bar{P}$$

After k cycles, the photon's energy is

$$\varepsilon_k = \varepsilon_i \bar{g}^k$$
 or $\ln \frac{\varepsilon_k}{\varepsilon_i} = k \ln \bar{g}$

Generation of power-law shaped spectra at high energies

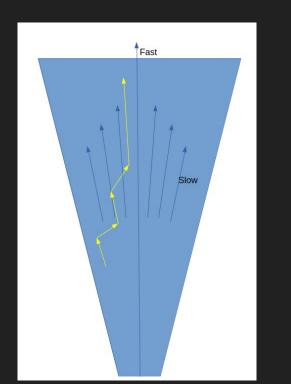
$$\ln \frac{N}{N_0} = \ln \frac{\varepsilon_k}{\varepsilon_i} \cdot \frac{\ln \bar{P}}{\ln \bar{g}}$$
$$\frac{N}{N_0} = \left(\frac{\varepsilon_k}{\varepsilon_i}\right)^{\beta'}$$
$$\beta' = \frac{\ln \bar{P}}{\ln \bar{g}}$$

$$\blacktriangleright \quad \frac{dN}{d\varepsilon_1} = N_0 \beta' \varepsilon_1^{\beta' - 1}$$

$$\beta = \beta' - 1 = \frac{\ln \bar{P}}{\ln \bar{g}} - 1$$

We supply P and g for calculating the spectral slopes at high energies.

2. Numerical simulations



Numerical Code:

- The numerical simulations are based upon Monte Carlo method
- We inject around 6 million photons deep inside the jet. Each of the photon goes through multiple scattering within the shear layers before it escapes.
- The escaped photons' population is distributed then into bins of observing angles as well as energies
- These binned photons produce the observed light curves as well as the spectrum for the given parameters

Results: Simulated spectrum (On axis observer)

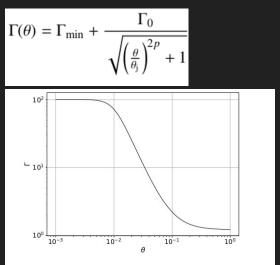
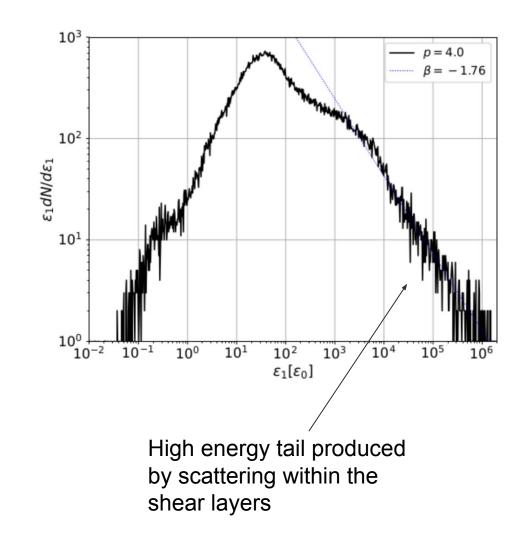
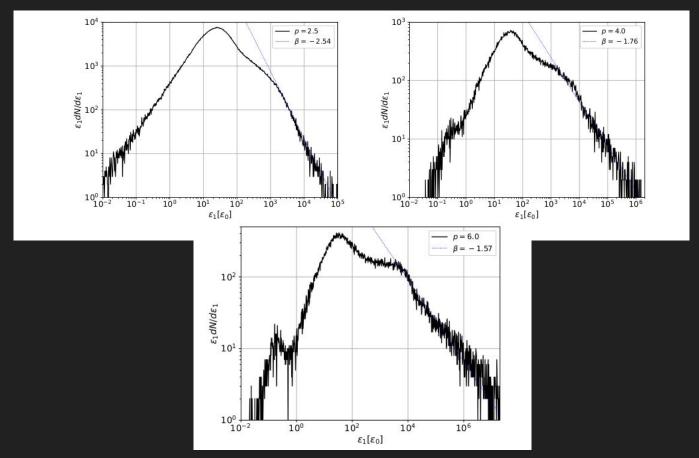


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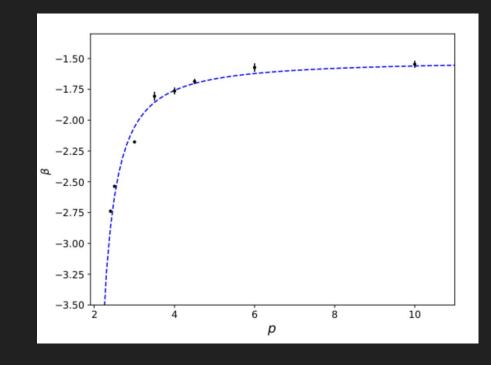
Simulated Spectra : Continued



Results : Analytic slopes and comparison with Monte Carlo Simulations

$$\beta = \beta' - 1 = \frac{\ln \bar{P}}{\ln \bar{g}} - 1$$

- Asymptotic slopes reach at -1.5 as p -> ∞
- For small and finite values of p (<2), the spectra vanishes due to geometric expansion of the jet
- The obtained range of β is compatible with the observed values



There are several other possible cases where such mechanism takes place

- A fast rotating star and its layers at the outskirts
- Relativistic turbulent plasma where there is a sharp velocity gradient in the turbulent eddies
- Sharply accelerating or decelerating relativistic jets/winds
- Corotating atmosphere of neutron stars (*Hanawa, ApJ, 1991*)

Summary : Conclusions and Significance

"Photons' Scattering in Relativistic Plasma with Velocity Shear: Generation of High Energy Power-law Spectra" Vyas and Pe'er, ApJL 943-1, L3, 7 Such power laws are capable of explaining high energy tails observed in GRB prompt phase spectra where β ranges between -3 and -1.5. We found the theoretical range $-\infty$ to -1.5.

Inversely, we show that using the observed values of β , we can directly constraint the jet structure of the these bursts.

The mechanism is important for emission from other such objects like AGNs or X ray binary jets, fast spinning compact stars, accretion discs with velocity gradient etc. We will explore these possibilities in future.