



A 4th-order accurate finite volume method for ideal and resistive classical and special relativistic MHD in the PLUTO code



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ABSTRACT

We present a novel implementation of a genuinely **4th-order accurate finite volume scheme** for multidimensional classical and special relativistic magnetohydrodynamics (MHD) in both ideal [1] and resistive [2] regimes based on the constrained transport (CT) formalism in the PLUTO code [3]. Our scheme is rooted over the method originally proposed by McCorquodale and Colella [4] but introducing several novel aspects when compared to its predecessors, yielding a **more efficient computational tool**. Among the most relevant ones, our scheme exploits **pointwise to pointwise reconstructions** (rather than one-dimensional finite volume ones), employs all the generalized upwind EMF averaging techniques of the UCT method of Mignone and Del Zanna [5], with the addition of a **new relativistic UCT-GFORCE** average, and **ensures robustness** through sophisticated limiting strategies that include both a discontinuity detector and an order reduction procedure. Such method has so far produced results in numerical simulations **that are unfeasible with traditional low order schemes**.

THE FINITE VOLUME FORMALISM IN (RELATIVISTIC) MHD

Our method is targeting the equations of ideal and resistive (R)MHD in Cartesian geometry.

1) **Hyperbolic sub-system** for the evolution of the conservative flow variables:

$$\frac{\partial \mathbf{U}}{\partial t} + \nabla \cdot \mathbf{F} = \mathbf{0},$$

$$\mathbf{U} = \begin{pmatrix} \rho \\ \rho \vec{v} \\ \mathcal{E} \end{pmatrix}, \mathbf{F} = \begin{pmatrix} \rho \vec{v} \\ \rho \vec{v} \vec{v} - \vec{B} \vec{B} + \mathbf{I} p_t \\ (\mathcal{E} + p_t) \vec{v} - (\vec{v} \cdot \vec{B}) \vec{B} \end{pmatrix}^T \quad \mathbf{U} = \begin{pmatrix} D \\ \vec{m} \\ \mathcal{E} \end{pmatrix}, \mathbf{F} = \begin{pmatrix} D \vec{v} \\ w_t \gamma^2 \vec{v} \vec{v} - \vec{B} \vec{B} - \vec{E} \vec{E} + \mathbf{I} p_t \\ \vec{m} \end{pmatrix}^T$$

MHD: with ρ density, \vec{v} fluid velocity, \vec{B} magnetic field, p_t total pressure, and \mathcal{E} energy density.

(Res)RMHD: with D relativistic mass density, \vec{m} momentum density, p_t total pressure, and \mathcal{E} energy density.

2) **Induction equation** for the evolution of the magnetic field:

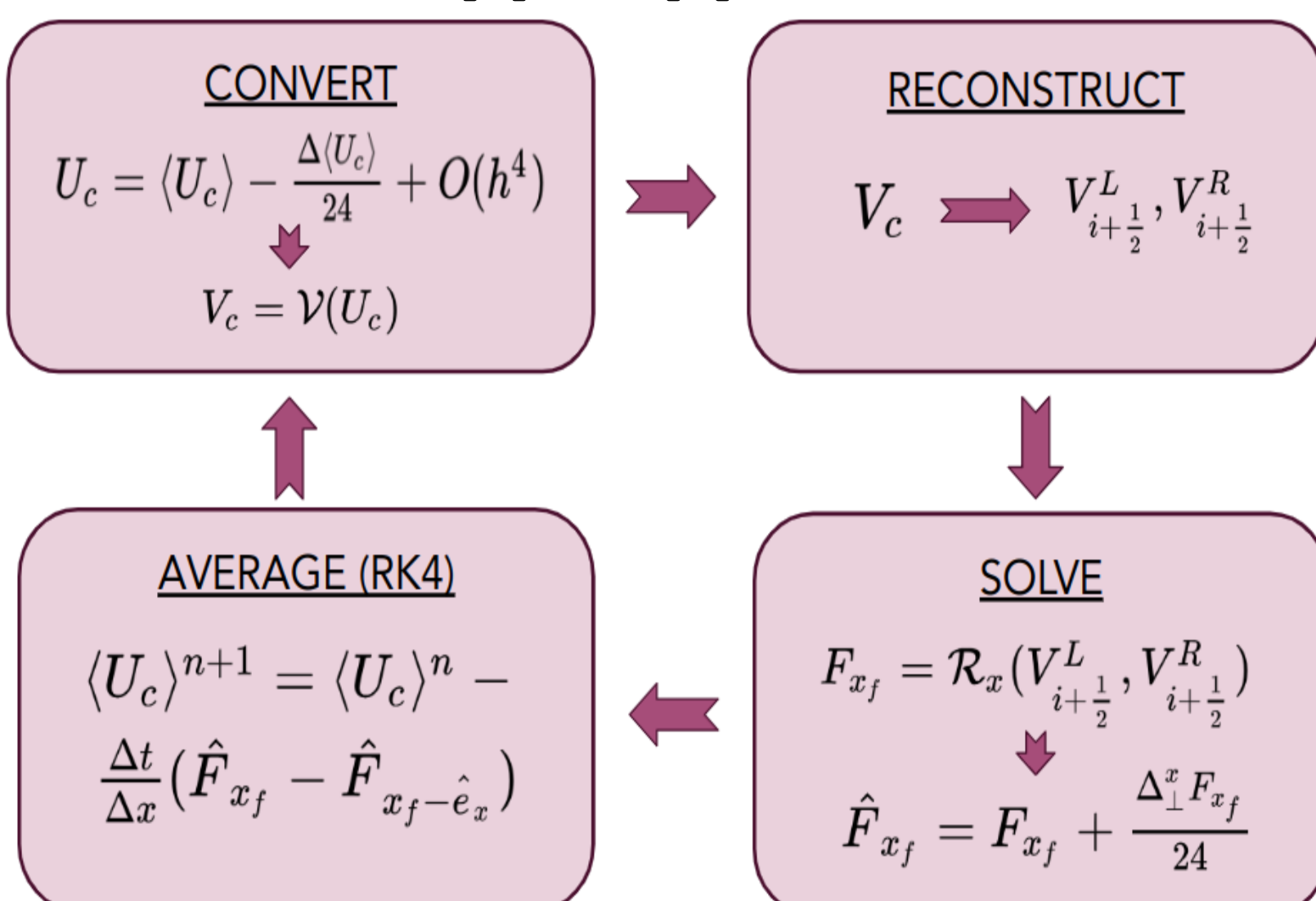
$$\frac{\partial \vec{B}}{\partial t} + c \nabla \times \vec{E} = \mathbf{0}.$$

ended with the **solenoidal condition** $\nabla \cdot \vec{B} = 0$ satisfied at $t = 0$ and preserved $\forall t > 0$.

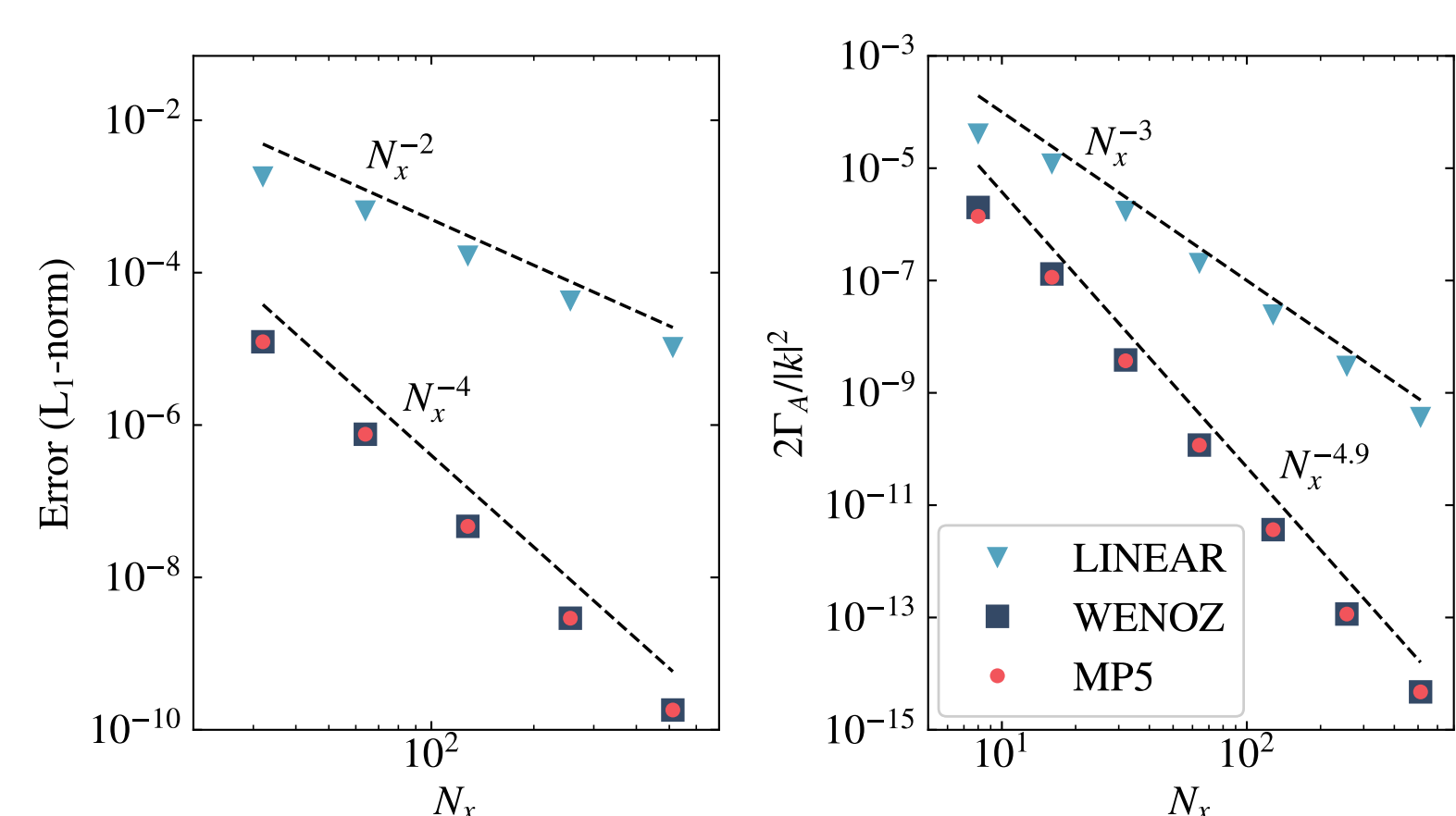
THE 4th-ORDER NUMERICAL SCHEME (IN BRIEF)

High order Godunov-type High-Resolution-Shock-Capturing scheme based on Reconstruct-Solve-Average strategy (for the details of the scheme see [1] and [2]):

- Conversion step** using Laplacian operators Δ [4] \Rightarrow obtain point values at cell centers V_c .
- Pointwise Reconstructions** [1] reconstruct V_c at interfaces without having to average $V_c \Rightarrow \langle V_c \rangle$.
- Limiter** [1] introduction of 2 types of discontinuity detectors [8],[9] and an order reduction procedure.
- Upwind Constrained Transport** + all the generalized upwind EMF averages [5] + punctual reconstructions.



THE 3D CIRCULARLY POLARIZED IDEAL ALFVÉN WAVES



- Convergence rates match the expected order \Rightarrow **improved accuracy**.
- For fixed accuracy $\epsilon_p = CN^{-p}$ ($p = 2, 4$) 2nd- and 4th-order scheme will achieve $\epsilon_2 \sim \epsilon_4$ with $N_4 \sim \sqrt{N_2}$ grid points.
- Net gain at fixed accuracy estimated: ϵ_2 at $N_2^3 = 512^3 \simeq \epsilon_4$ at $N_4^3 = 34^3$ + computing time reduced by a factor $10^4 \Rightarrow$ **saving of computational time**.

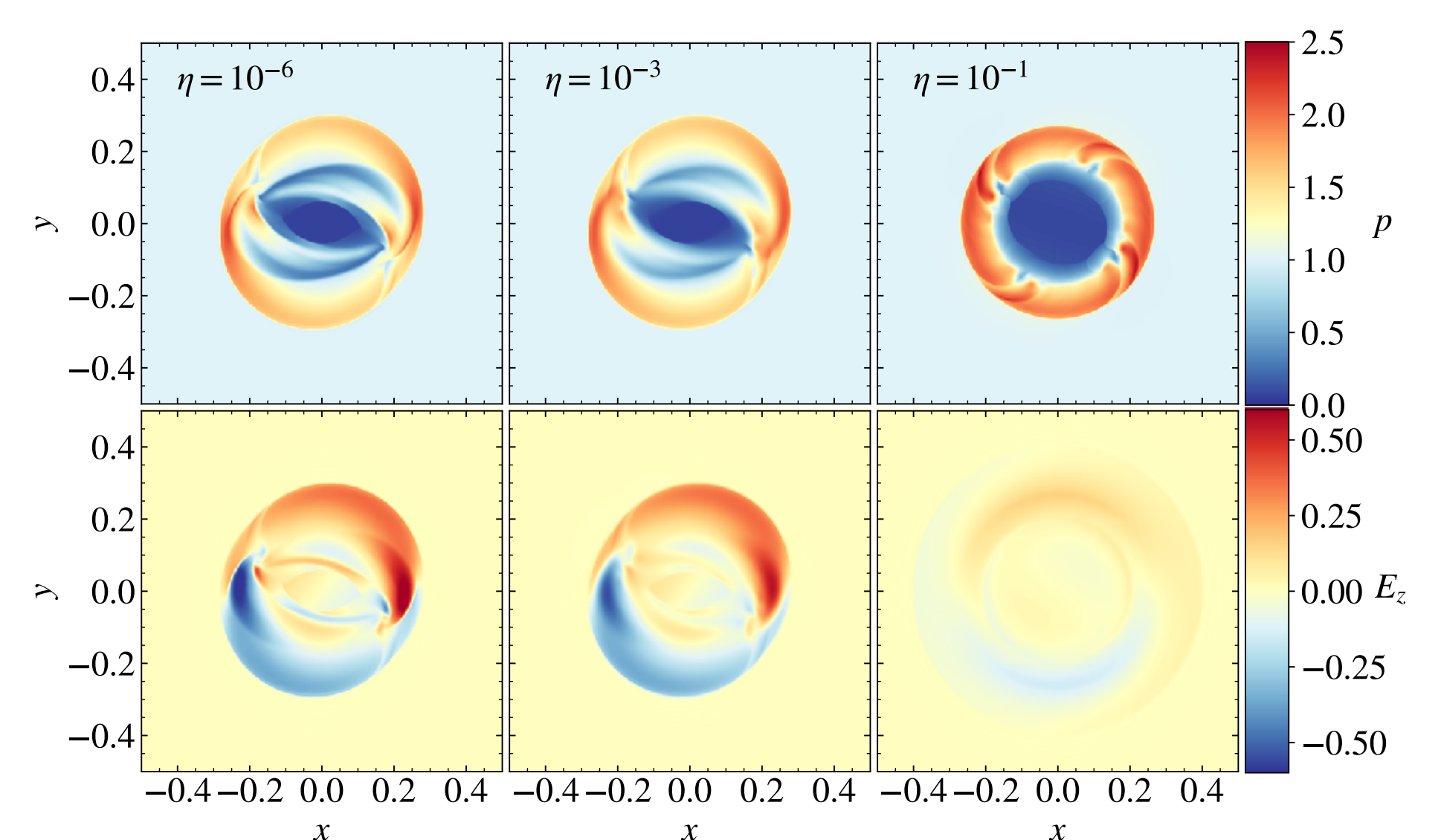
- More accurate scheme \Rightarrow **reduced numerical dissipation**.
- Overall decay rate of $\sim N^{-4.9}$ for the 4th-order scheme versus only $\sim N^{-3.0}$.

RESULTS, CURRENT AND FUTURE DEVELOPMENT

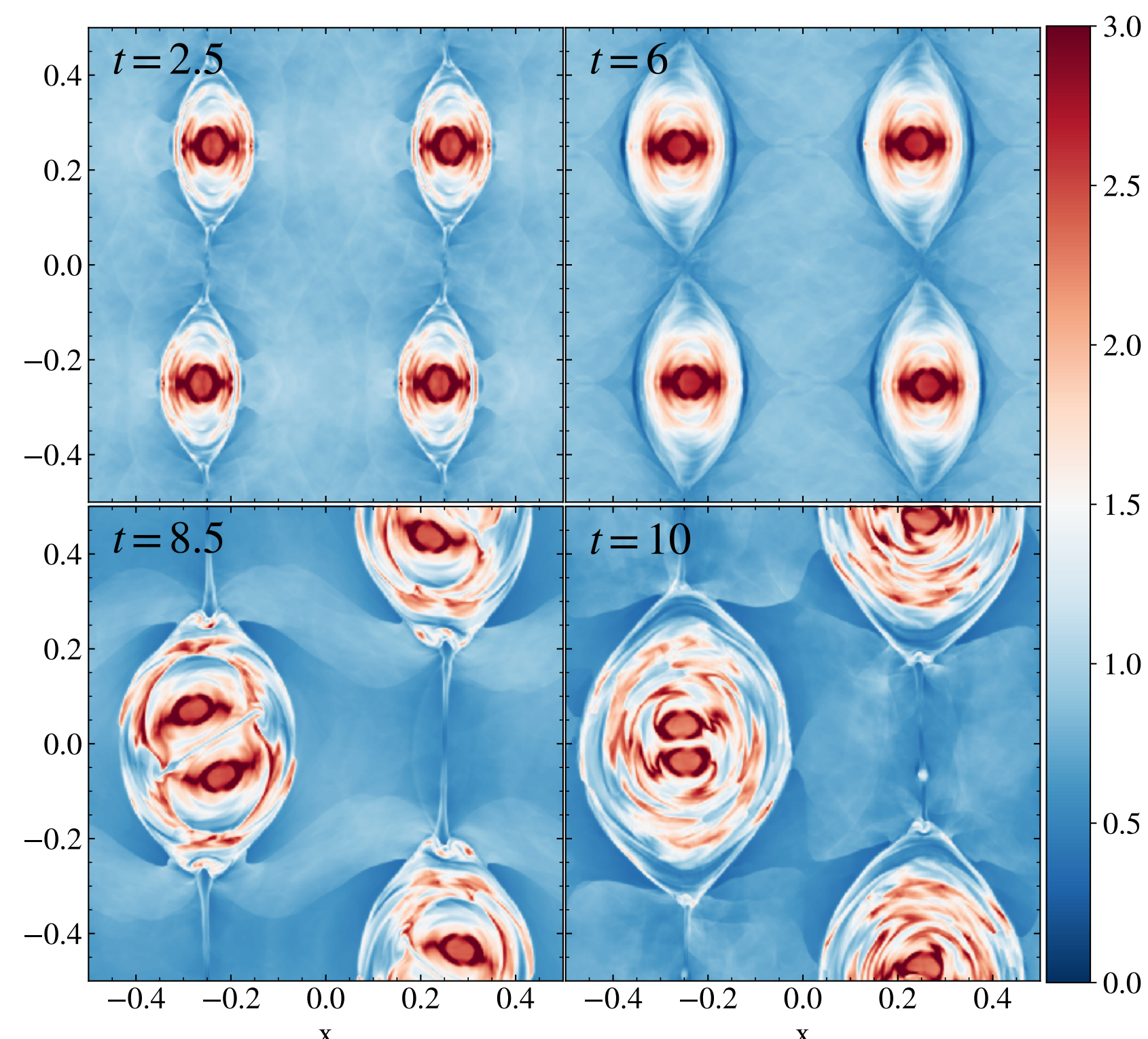
- Genuinely 4th-order accurate finite volume scheme for both ideal and resistive (R)MHD accounting several innovative aspects that yield an accurate, robust and efficient computational tool;
- Introduction of pointwise to pointwise reconstruction operations that ease up the structure of the scheme;
- Robustness assured by a limiter that allows the 4th-order scheme to carry out extremely severe tests yielding unprecedented results;
- Generalization of the UCT-GFORCE average to relativistic MHD;

THE RESRMHD ROTOR

- High density disk rotating at relativistic speed ($\omega = 8.5$) simulated with increasing resistivity: $\eta = 10^{-6}, 10^{-3}, 10^{-1}$.
- The complex pattern of shocks and torsional Alfvén waves is correctly reproduced by the new scheme.
- Magnetic braking slows down the rotor and **physical resistivity** dissipates the electric field.



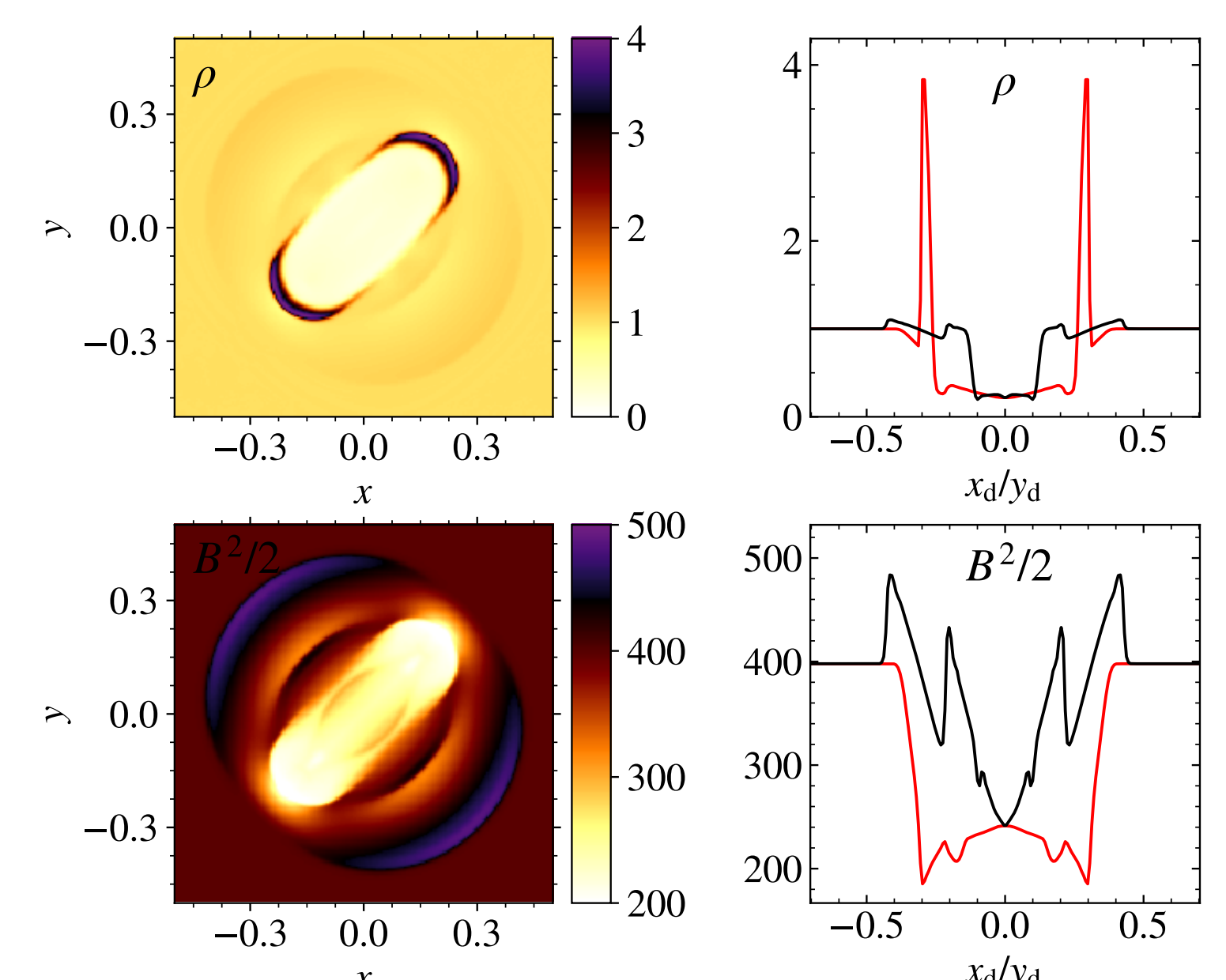
THE RELATIVISTIC CURRENT SHEET



- Upwind Constrained Transport method using all the EMF averages of [5] to ensure $\nabla \cdot \vec{B} = 0$.
- Pointwise reconstructions reduce the number of transverse ghost zones to reconstruct the electric field \Rightarrow **more efficient CT algorithm**.
- Extended the MHD **UCT-GFORCE** average to RMHD \Rightarrow accurate and robust solver.
- This ideal relativistic current sheet reproduces the one by [6], but with a more accurate and less dissipative solver (GFORCE rather than linear solver).

THE 3D MHD BLAST WAVE

- Severe blast configuration [7] with $B_0 = 100/\sqrt{4\pi}$, $p_{out} = 1$ and $p_{in} = 10^3$.
- So far only feasible with 2nd-order schemes.
- Limiter acts in two stages of the algorithm:
 - during the conversion $\langle U_c \rangle \Rightarrow U_c$ for **reverting to 2nd-order** with Jameson's shock sensor [8] or by a more sensitive limiter based on higher order derivatives' ratio [9];
 - during reconstruction at interfaces by means of an **order reduction** procedure when oscillations are generated.



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