

High energy neutrino emission from a radiatively inefficient accretion flow based on a GRMHD simulation

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Introduction

• High energy neutrino :

- ✓ Emitted via pp and $p\gamma$ collision processes of accelerated protons
- ✓ Trajectories are not affected by B-field attributed to the neutral charge
- They can be a "smoking gun" of cosmic-ray (CR) acceleration

• Sources of IceCube neutrinos: uncertain

- ✓ Active Galactic Nuclei (AGN)
- ✓ Galaxy Clusters
- ✓ Starburst Galaxies
- ✓ Low Luminosity Gamma-ray Bursts

• For more quantitative studies, the theoretical works based on the global structure of the plasma, i.e., models beyond the I-zone approximation, will be important to be explored.

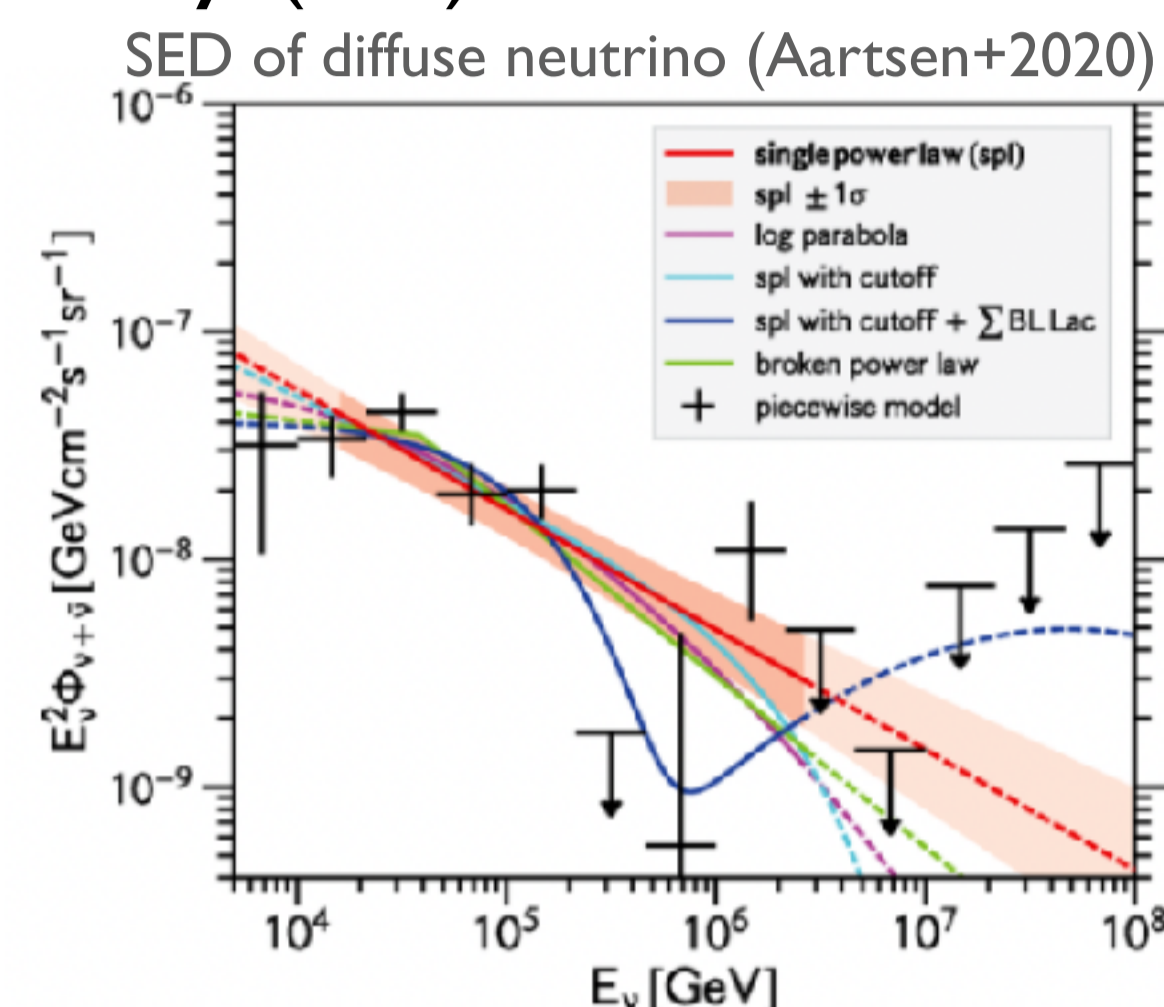
ex) AGNs

- ✓ Global structure and states of the accretion flow (incl. magnetization)
- ✓ Turbulence in kinetic scale
- ✓ Black hole spin

For a pioneering work of neutrino emission model with I-zone approximation, see Kimura et al. (2015), etc.

In this work, we have developed ν -RAIKOU code, which calculates the CR proton (CRp) acceleration and neutrino emission using 3D General Relativistic MHD (GRMHD) data.

We have studied the effects of the **global structure** of accretion flows on the resultant neutrino SEDs.

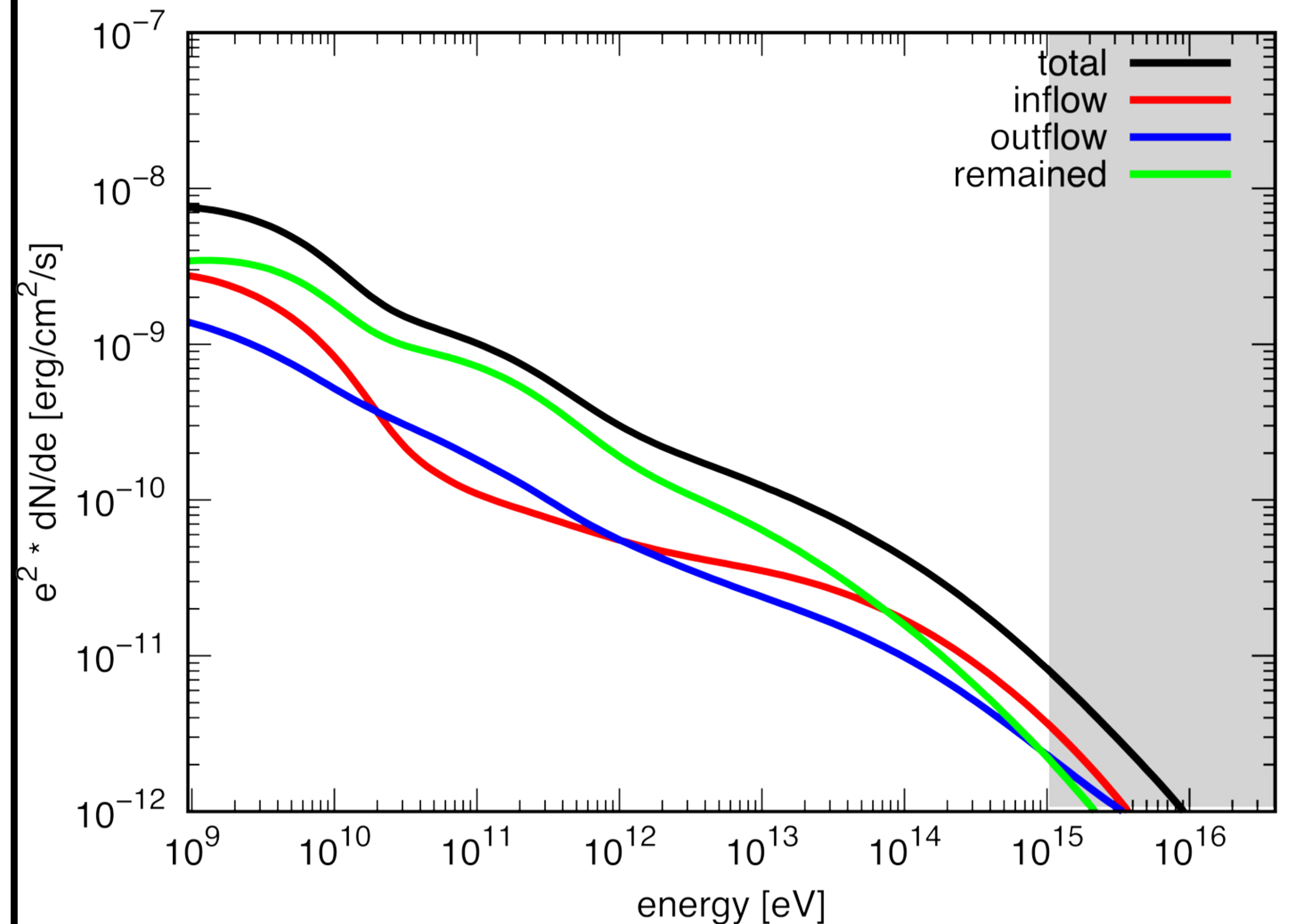


- ✓ CR acceleration
- ✓ Neutrino SEDs

Result

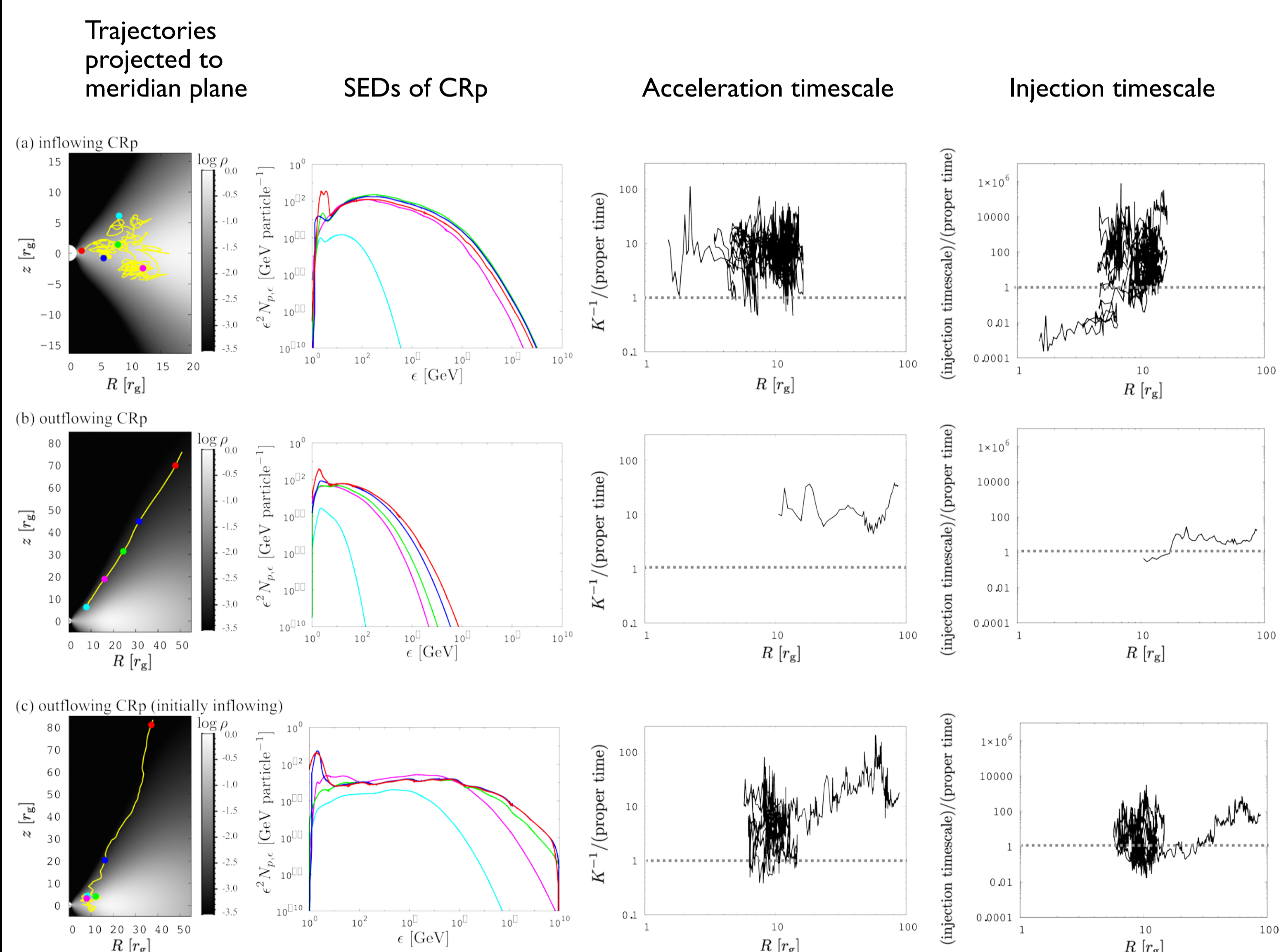
• Time averaged neutrino SED

- ✓ SEDs flatter than those of I-zone models
- ✓ Neutrinos originated from CRp finally captured by black hole (inflow) contribute as much as those from outflowed CRp.
- ✓ These moderately flat SED may explain the origin of diffuse SEDs.



• Trajectories, SEDs and acceleration timescale of CRp:

The CRp with various trajectories and acceleration forms the flatter neutrino SEDs.



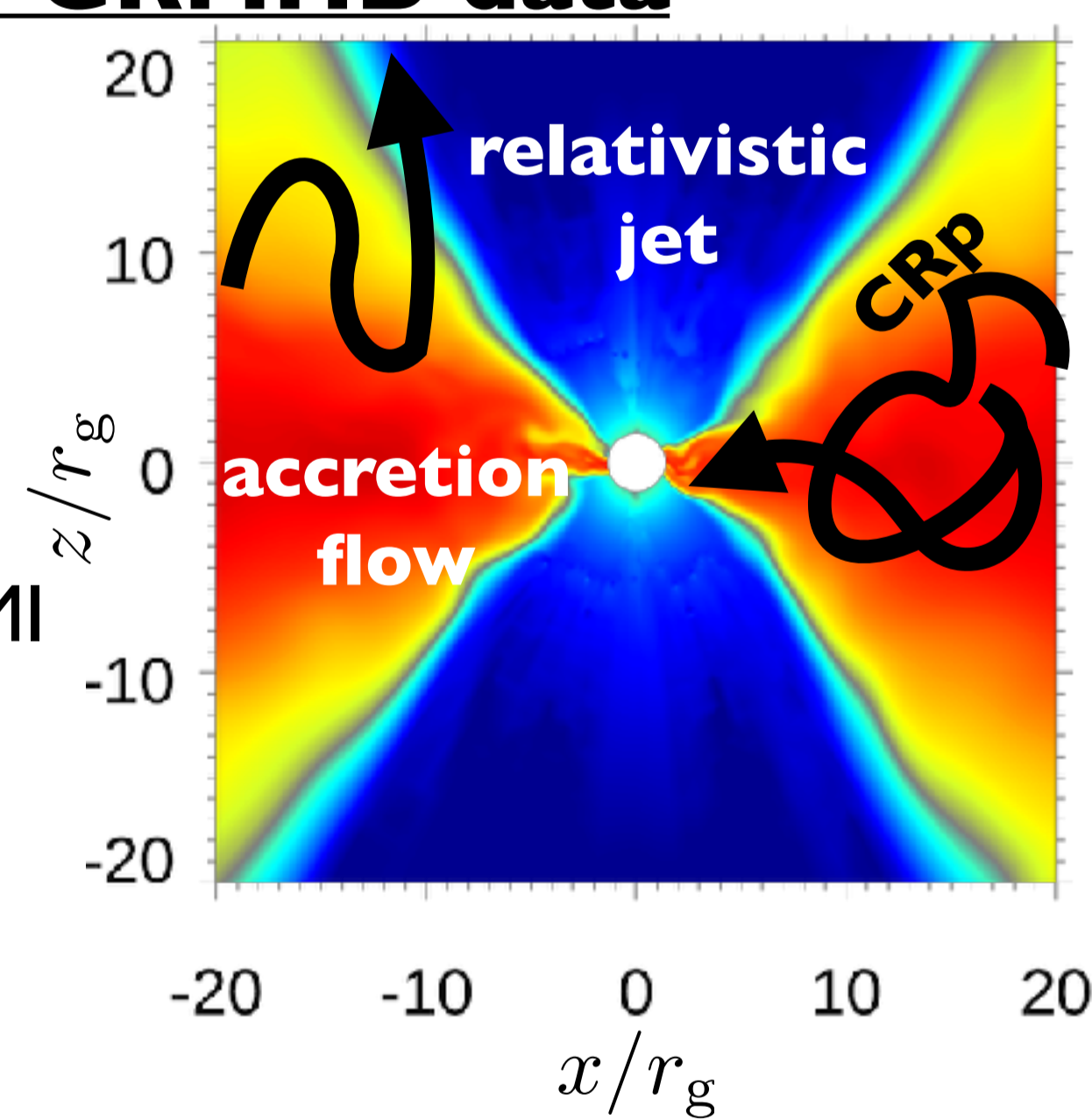
Method

(1) Trajectories of CRp using 3D GRMHD data

• The trajectories along the stream lines

We are interested in the acceleration up to \sim PeV in this work where the gyro radius is less than the mesh size.

• GRMHD data (Kawashima+2023) in semi-MAD state (intermediately magnetized accretion flow). GR(R)MHD code UWABAMI (Takahashi + 2016)



(2) Time evolution of CRp SEDs

Fokker-Planck eq.

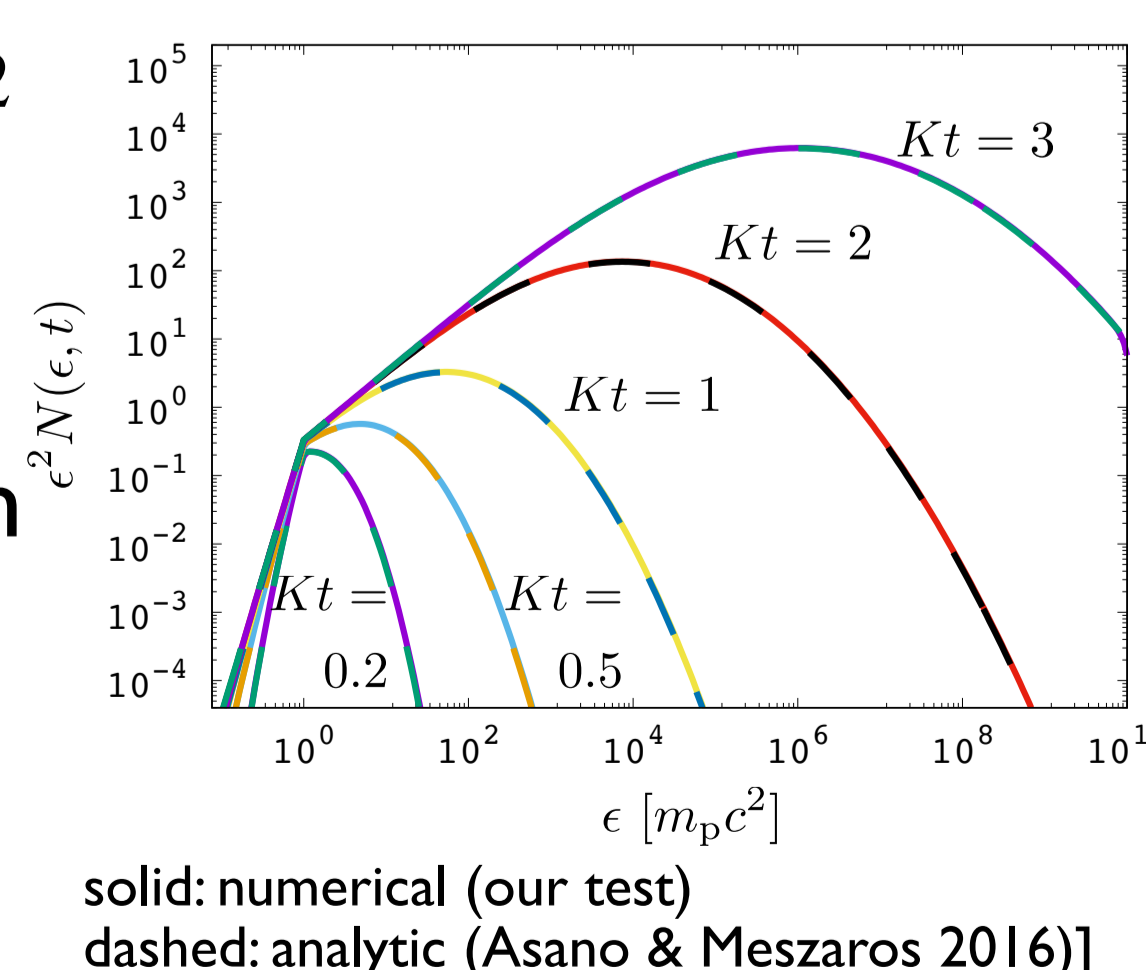
$$\frac{\partial N(\epsilon', t')}{\partial t'} = \frac{\partial}{\partial \epsilon'} \left[D(\epsilon') \frac{\partial N(\epsilon', t')}{\partial \epsilon'} \right] - \frac{\partial}{\partial \epsilon'} \left[\frac{2D(\epsilon')}{\epsilon'} N(\epsilon', t') \right] + \dot{N}_{inj}(\epsilon', t')$$

Solved with a method with Green function (Becker + 2006).

• energy diffusion: turbulent acceleration w/ hard sphere approx., i.e., $D(\epsilon') = K\epsilon'^2$ (ϵ' : CRp energy in the fluid-rest frame)

• injection: CRp of $\epsilon' = 2m_p c^2$ w/ higher injection rate at highly magnetized region (motivated by magnetic reconnection).

• Effects of Compression/Rarefaction are also included.



(3) High-energy neutrino SED

- pp collision between CRp and MHD (thermal) proton.
- Neutrinos w/ approx. formula of pion SED (Kelner + 2006)
- The effects of Gravitational redshift are included.

Summary and Prospects

• We have computed neutrino SEDs of global accretion flow based on GRMHD model, with developing a new code calculating CRp accelerations and neutrino emission using GRMHD simulation data (ν -RAIKOU code).

• It is found that neutrino SEDs flatter than those of I-zone model appears due to the superposition of emission from various CRp attributed to the global structure of the accretion flow.

• The moderately flat SED may explain the origin of the diffuse neutrino observed by IceCube.

• We will add the effects of the $p\gamma$ processes, combining the ν -RAIKOU with general relativistic multiwavelength radiative transfer code RAIKOU (Kawashima+2023)